



# The effects of the surface inflows on quenching of solar cycles



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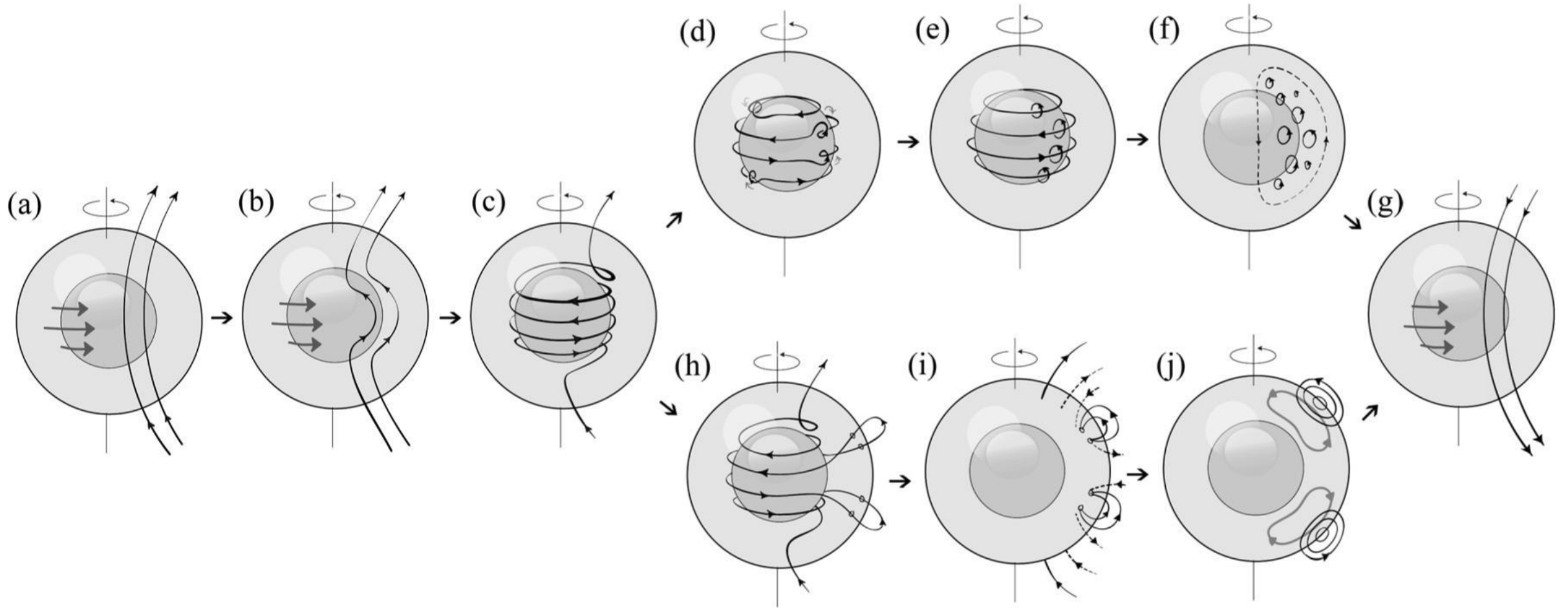
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# Introduction



Sanchez et. al. 2014, An. Acad. Bras. Ciênc. 86 (1), 11–26.

# Motivations

- Understand the effects of surface inflows on the solar dynamo.
- Investigate the quenching parameters in Surface Flux Transport models (SFT).
- Explore the effects of surface inflows on the SFT's in the presence of quenching mechanisms.

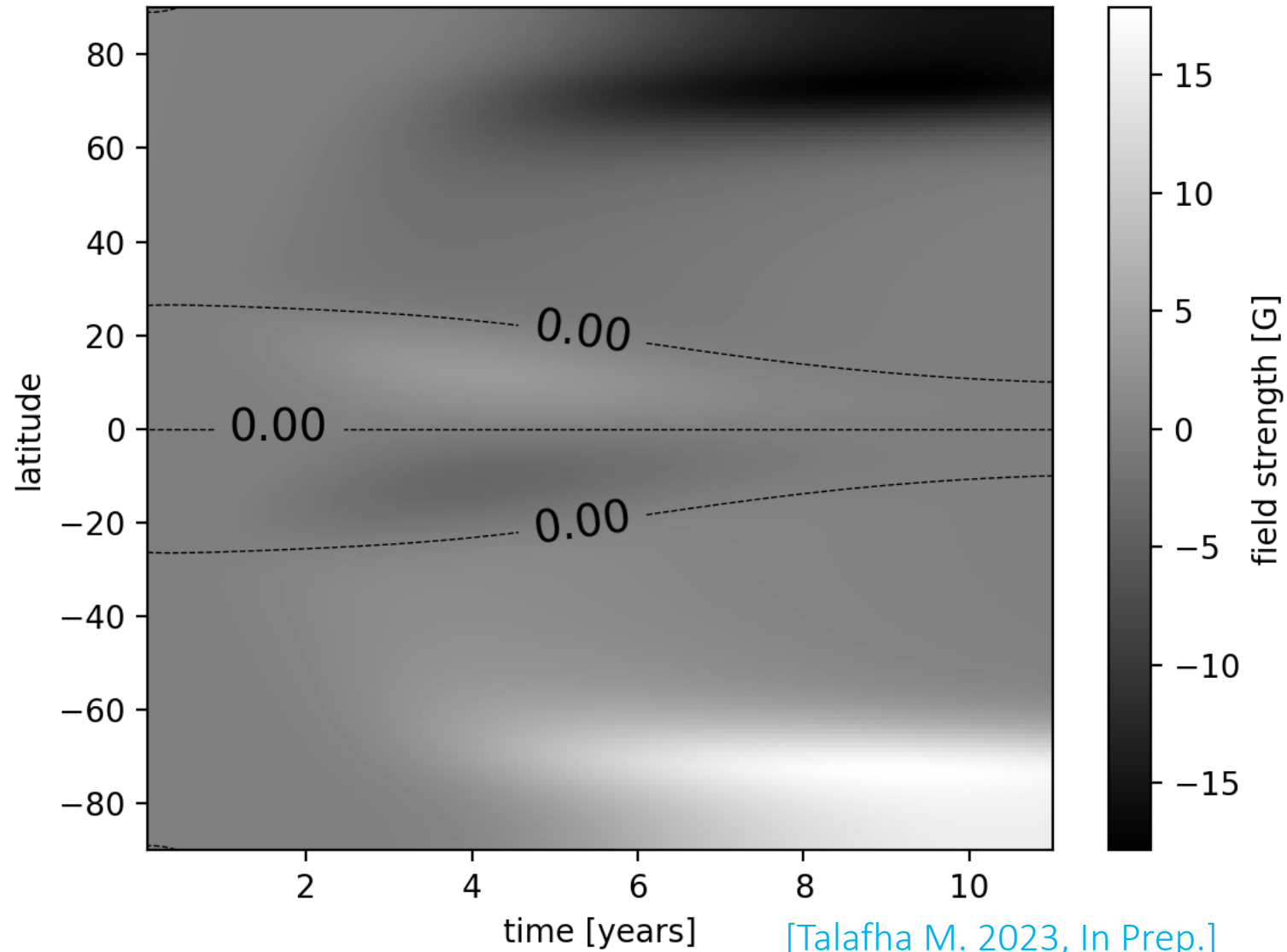
# Surface Flux Transportation Model (SFT)

$$\frac{\partial B}{\partial t} = \nabla \times (u \times B - \eta \nabla \times B)$$

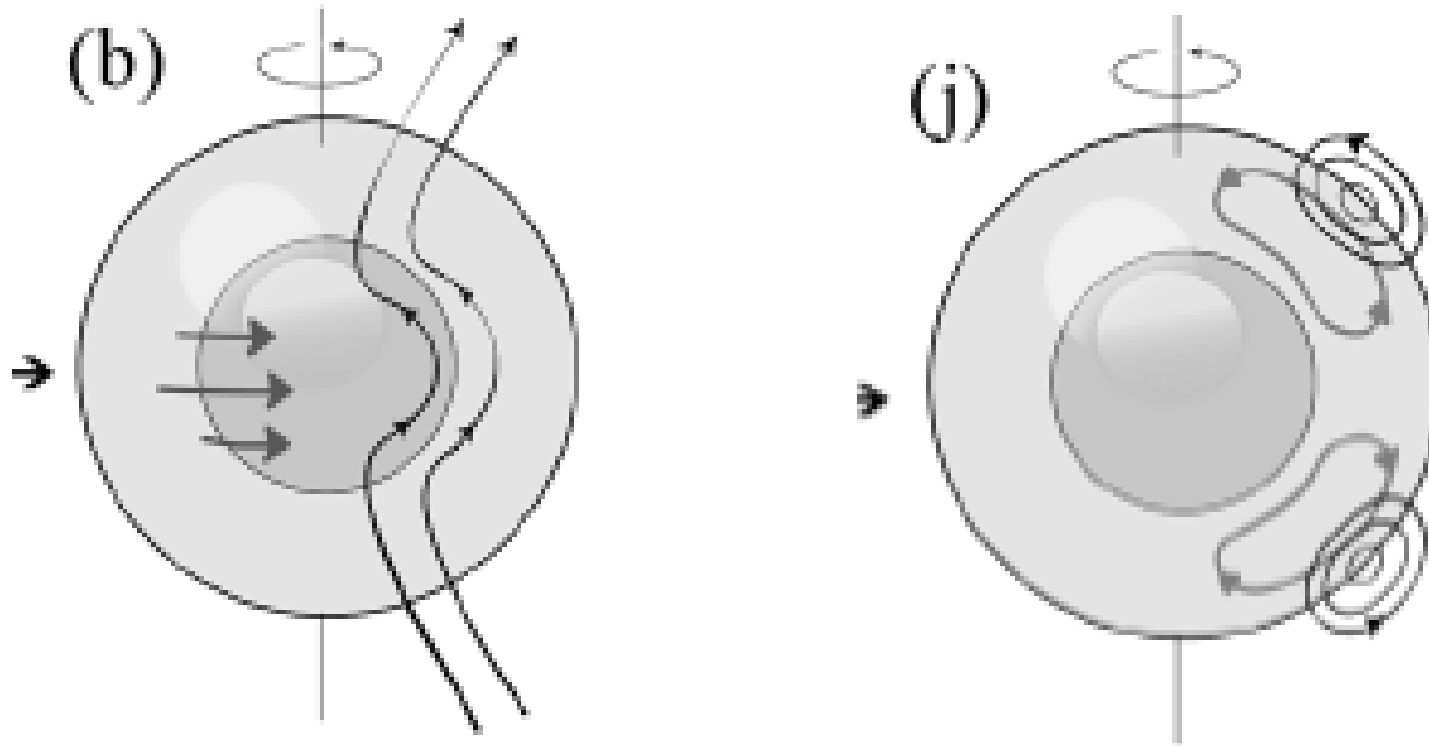
- considering the magnetic field to be predominantly radial on global scales solving only the r-component, after enforcing the null divergence condition throughout:

$$\begin{aligned} \frac{\partial B_R}{\partial t} = & -\frac{1}{R \sin \theta} \frac{\partial}{\partial \theta} [u_\theta(R, \theta) B_R \sin \theta] - \Omega(R, \theta) \frac{\partial B_R}{\partial \phi} \\ & + \frac{\eta_R}{R^2} \left[ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial B_R}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 B_R}{\partial \phi^2} \right] - \frac{B_R}{\tau_R} + S_{BMR}(\theta, \phi, t) \end{aligned}$$

Flow 2,  $u_0 = 10.0$ ,  $\eta = 250.0$ ,  $\tau = 8.0$ , cycle=1

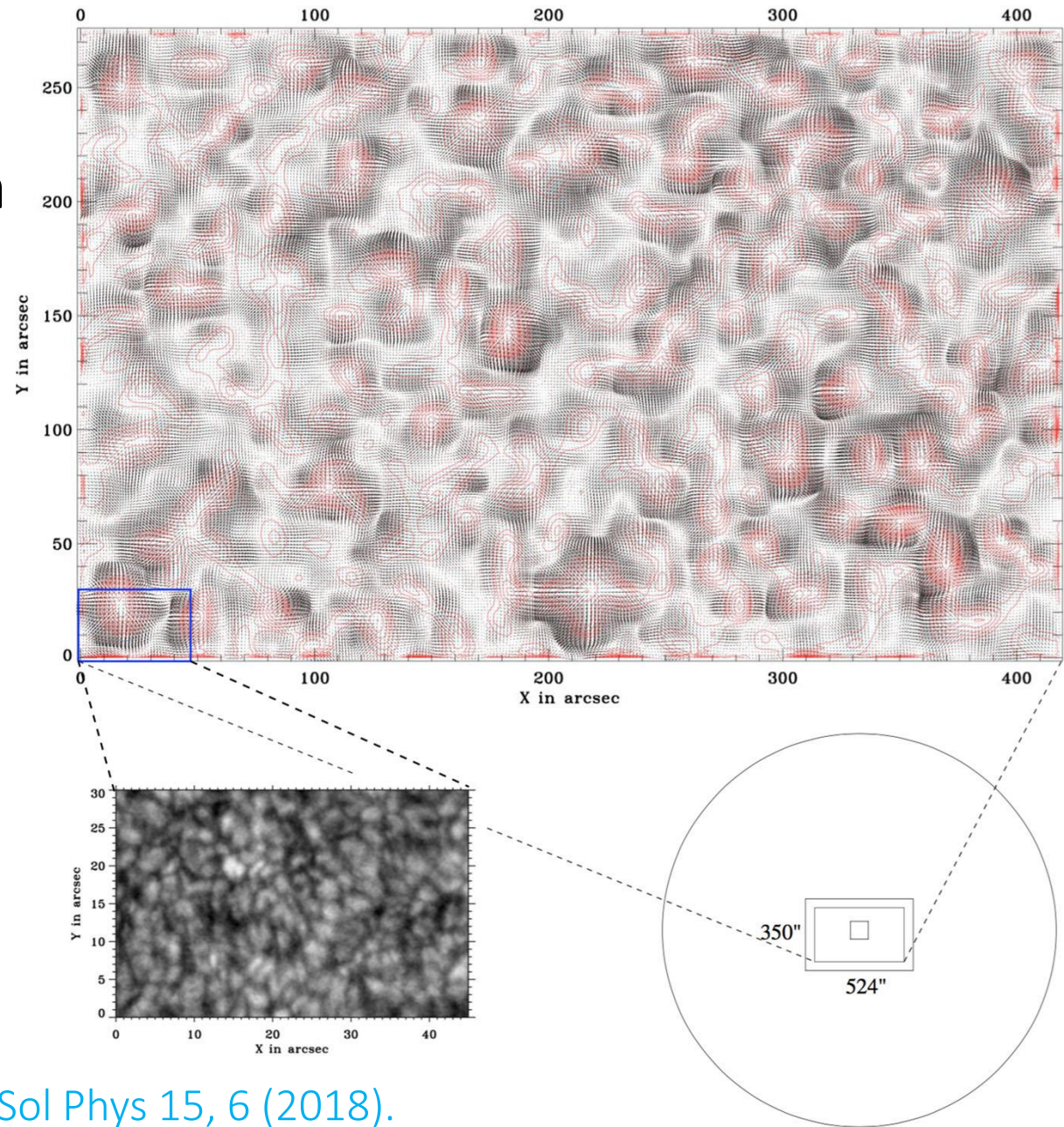


- Large scale flows: Differential rotation and Meridional flow



Sanchez et. al. 2014, An. Acad. Bras. Ciênc. 86 (1), 11–26.

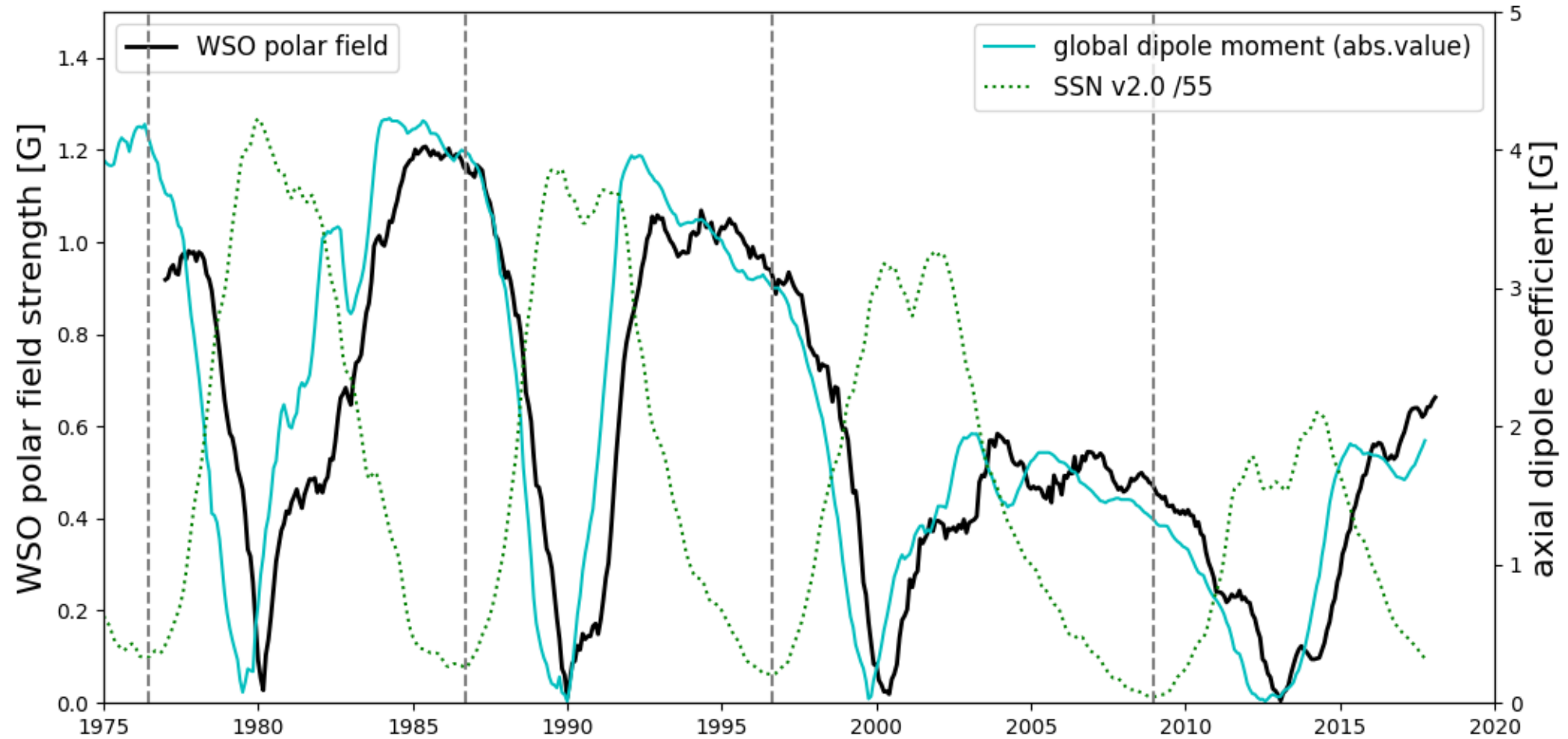
- Small scale flows: diffusion



Rincon, F., Rieutord, M.. [Living Rev Sol Phys 15, 6 \(2018\)](#).



# Constraints



Petrovay, K. & Talafha, M. 2019, A&A, 632, A87

# Computations

- Python script include (FTCS scheme)

1- **Source function:** represented by a pair of rings of opposite magnetic polarity.

$$\begin{aligned}
 S(\lambda, t) = & kA_m S_1(t) S_2(\lambda; \lambda_o(t) - \Delta\lambda(t), \delta\lambda) \\
 & - kA_m S_1(t) S_2(\lambda; \lambda_o(t) + \Delta\lambda(t), \delta\lambda) \\
 & + kA_m S_1(t) S_2(\lambda; -\lambda_o(t) - \Delta\lambda(t), \delta\lambda) \\
 & - kA_m S_1(t) S_2(\lambda; -\lambda_o(t) + \Delta\lambda(t), \delta\lambda)
 \end{aligned}$$

$k$  : is a factor alternating between even and odd cycles.

$A_m$  : is an arbitrary amplitude depending on the flow profile (0.015 for the sinusoidal case.)

$S_2(\lambda; \lambda_o(t), \delta\lambda)$  is the latitudinal profile represented by a Gaussian, migrating equatorward during the course of a cycle.

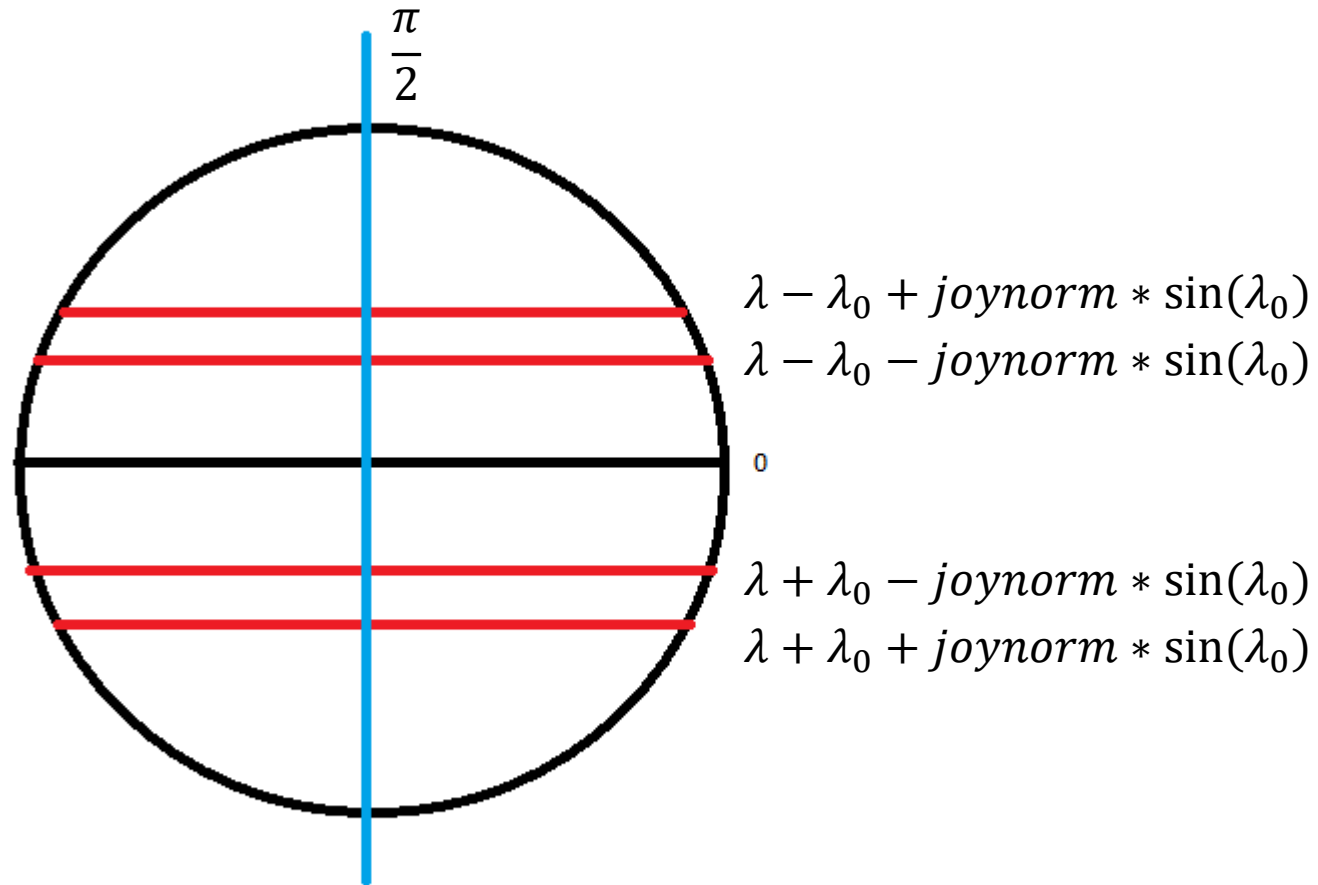
the time profile of solar activity in a typical cycle was determined by Hathaway et al. 1994, from the average of many cycles

Their trajectory during the course of a cycle is given by a quadratic fit from Jiang et al. 2011.

The latitudinal separation of the rings is a consequence of Joy's law.

$$\lambda_o[o] = 26.4 - \frac{34.2t}{P} + 16.1\left(\frac{t}{P}\right)^2 \quad 2\Delta\lambda = 0.5 \frac{\sin \lambda}{\sin 20^\circ} \quad S_1(t) = at_c^3 / [\exp\left(\frac{t_c^2}{b^2}\right) - c]$$

# Latitudes of pair of rings of opposite polarity



## The latitudinal profile ( $S_2$ ):

$$s_2(\lambda; \lambda_o, \delta\lambda) = \frac{\delta\lambda_o}{\delta\lambda} \exp\left[-\frac{(\lambda - \lambda_o)^2}{2\delta\lambda^2}\right]$$
$$\delta\lambda = \left[0.14 + 1.05\left(\frac{t}{P}\right) - 0.78\left(\frac{t}{P}\right)^2\right] \lambda_o$$
$$\delta\lambda_o = 6.26^\circ$$

[Jiang, et. al. 2011, A&A, 528, A82]

## Meridional flow profiles :

- (Van Ballegooijen et al. 1998); used by (Jiang et al. 2014).

$$u_c = \begin{cases} u_o \sin(\pi \times \lambda/90), & \text{if } |\lambda| < \lambda_o \\ 0 & , \text{otherwise} \end{cases}$$

- Based on the optimized parameters [Petrovay, K. & Talafha, M. 2019]
  - **Diffusivity:** values ( $\eta$ ) (250 and 600)  $\text{km}^2 \text{s}^{-1}$ .
  - **Source decay time:** values ( $\tau$ ) (8 - 100) yr.
  - **Meridional flow velocity:** values ( $u_o$ ) (11)  $\text{ms}^{-1}$ .

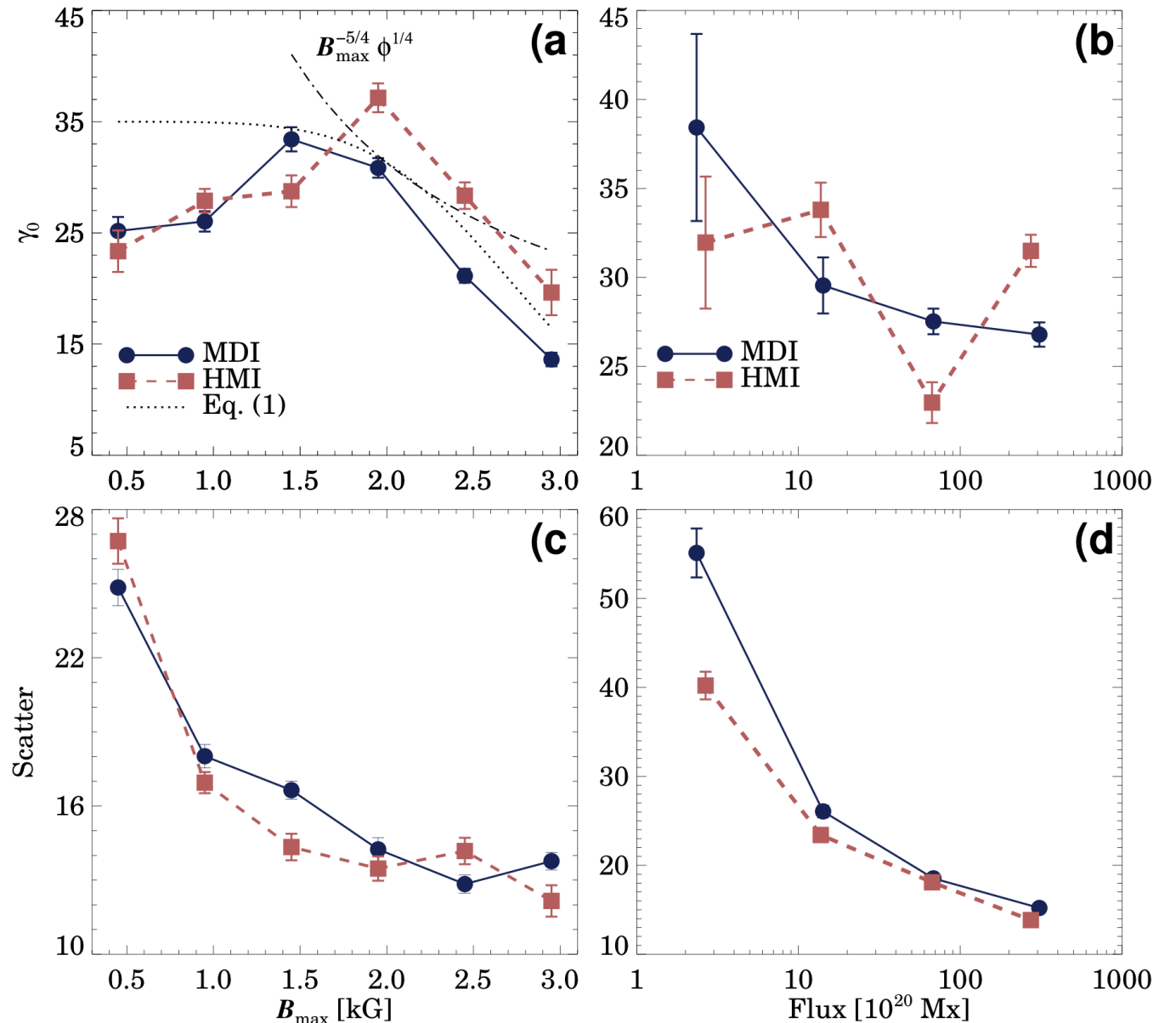
Parameter	NoQ	Different cases of LQ and TQ				
$b_{\text{lat}}$	0	1.5	2.4*	3.5	4.5	5.5
$b_{\text{joy}}$	0	0.05	0.1	0.15*	0.2	0.25

\* standard values [Jiang, et. al. 2011, A&A, 528, A82]

- For each combination of the parameters, we run the code solving SFT equation for 1000 solar cycles to produce a statistically meaningful sample of cycles.

# Nonlinearities in solar cycles

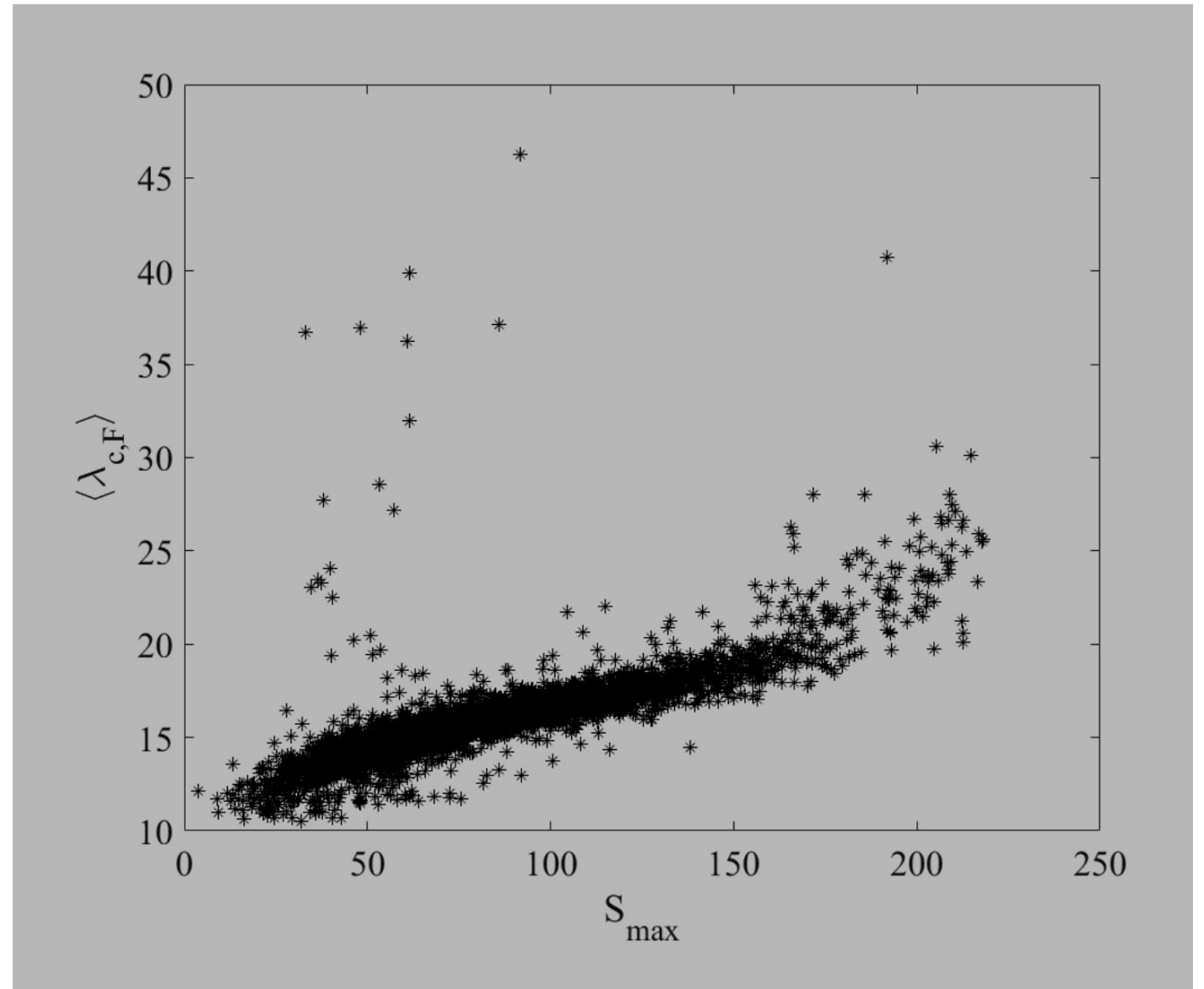
- Tilt quenching (TQ)



Jha, B. K., Karak, B. B., Mandal, S., & Banerjee, D. (2020). The Astrophysical Journal Letters, 889(1), L19.

# Nonlinearities in solar cycles

- Latitude quenching (LQ)
  - A positive correlation between the cycle amplitude and average emergence latitude of active regions



[Talafha, et. al. 2022, A&A, 660, A92]

## Tilt quenching (TQ)

$$\Delta\lambda = 1.5^\circ \sin\lambda_o \left( 1 - b_{joy} \frac{A_n - A_0}{A_0} \right)$$

[Talafha, et. al. 2022, A&A, 660, A92]

## Latitude quenching (LQ)

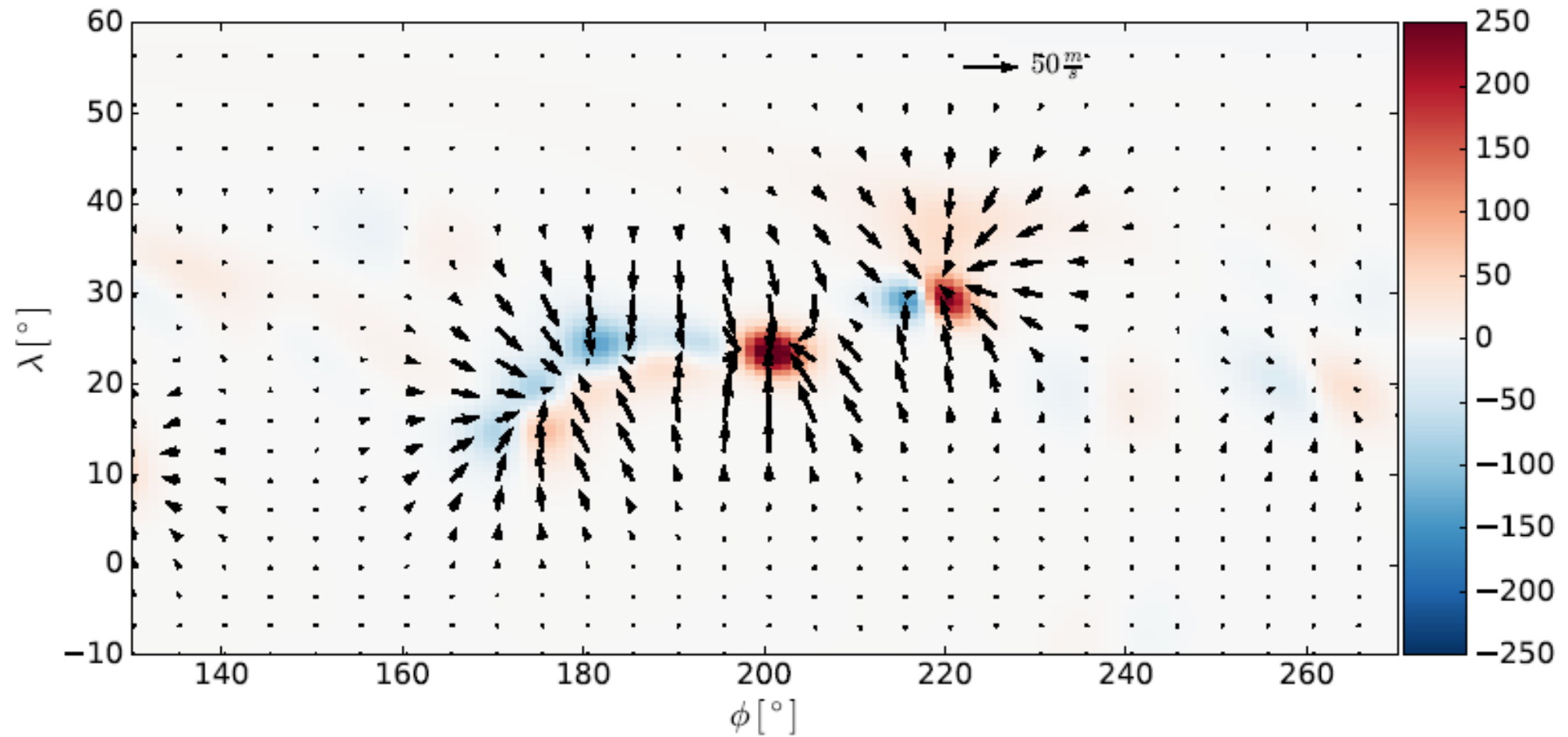
$$\lambda_o(t; i) = \left[ 26.4 - 34.2 \left( \frac{t}{P} \right) + 16.1 \left( \frac{t}{P} \right)^2 \right] \left( \frac{\lambda_i}{14.6^\circ} \right)$$

$$\lambda_i = 14.6^\circ + b_{lat} \frac{A_n - A_0}{A_0}$$

[Jiang, et. al. 2011, A&A, 528, A82]



# Surface inflows



Martin-Belda, D., & Cameron, R. H. (2017). *A&A*, 597, A21.

$$u_c = A \times \begin{cases} u_o \sin(\pi \times \lambda/\lambda_o) + \Delta v(\lambda, t), & \text{if } |\lambda| < \lambda_o \\ 0 & \text{, otherwise} \end{cases}$$

$$\Delta v(\lambda, t) = \begin{cases} v_o \sin(\lambda - \lambda_c/\Delta\lambda_v), & -180 < [\lambda - \lambda_c/\Delta\lambda_v] < 180 \\ 0 & \text{, otherwise} \end{cases}$$

[Jiang et. al. 2010. The Astrophysical Journal 717(1), 597]

Where:

$$\lambda_o = 75^\circ$$

$$\Delta\lambda_v = 15^\circ$$

$$v_o = 0 \text{ and } 5 \text{ m/s}$$

$$\lambda_c = 0, \pm 5, \pm 15, \pm 25$$

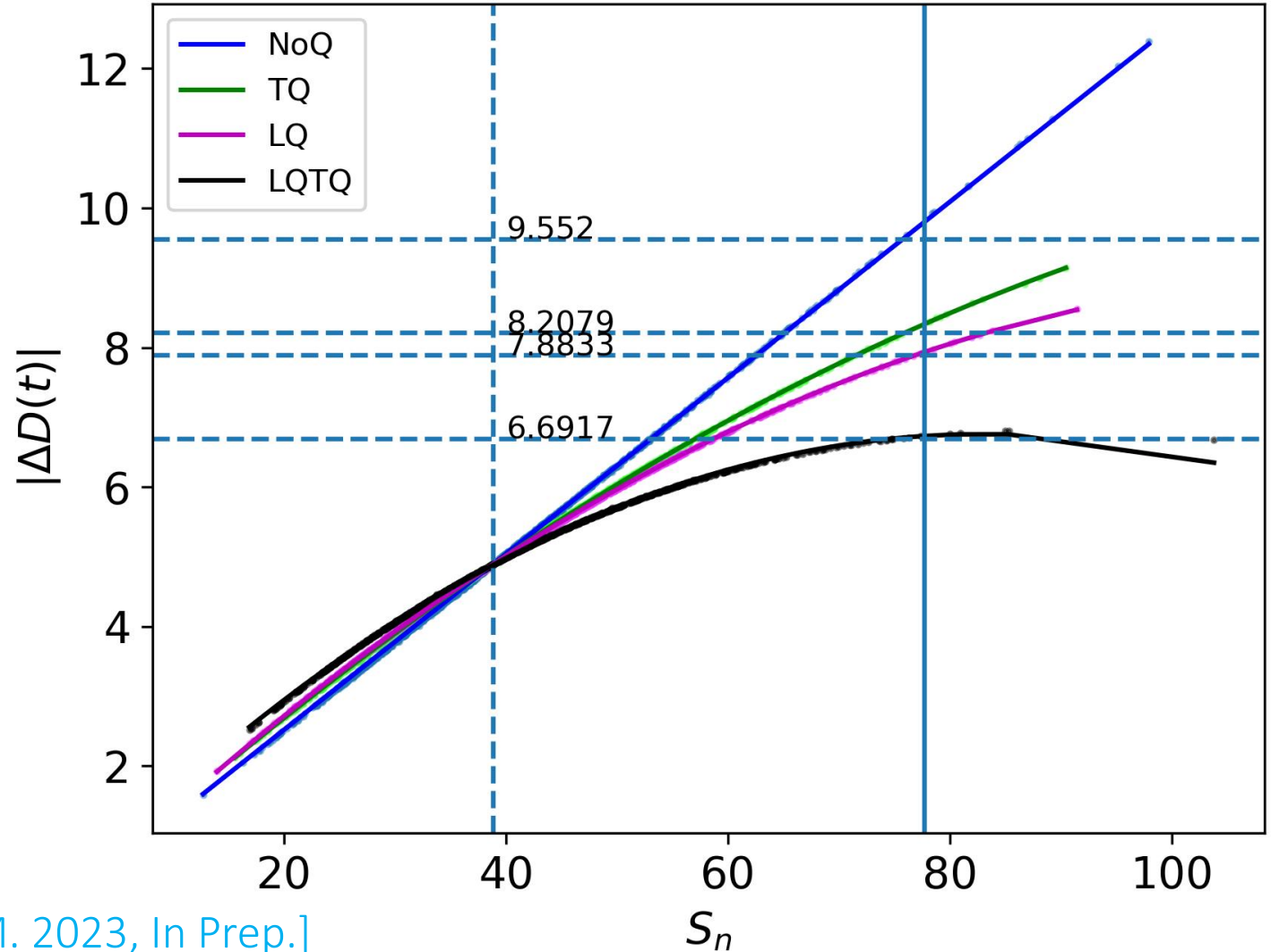
- The amplitude of the perturbed meridional flow (A) was scaled to be depending on the cycle amplitude as  $([A_n - A_0]/A_0)$
- $A_0$  is an arbitrary reference value corresponding to a “typical” or average cycle amplitude.

# Preliminary Results

- A linear dependence of both quenching parameters ( $b_{\text{joy}}$  and  $b_{\text{lat}}$ ) on the dipole moment of the solar cycle.
- The build up of the surface inflows in the cases of  $\lambda_c = \pm 5$  and  $\pm 25$  was very quick, it covered the entire surface by the first 20 solar cycles.

$$u_0 = 10.0, \eta = 250, \tau = 8, v_0 = 0, \lambda_c = 0$$

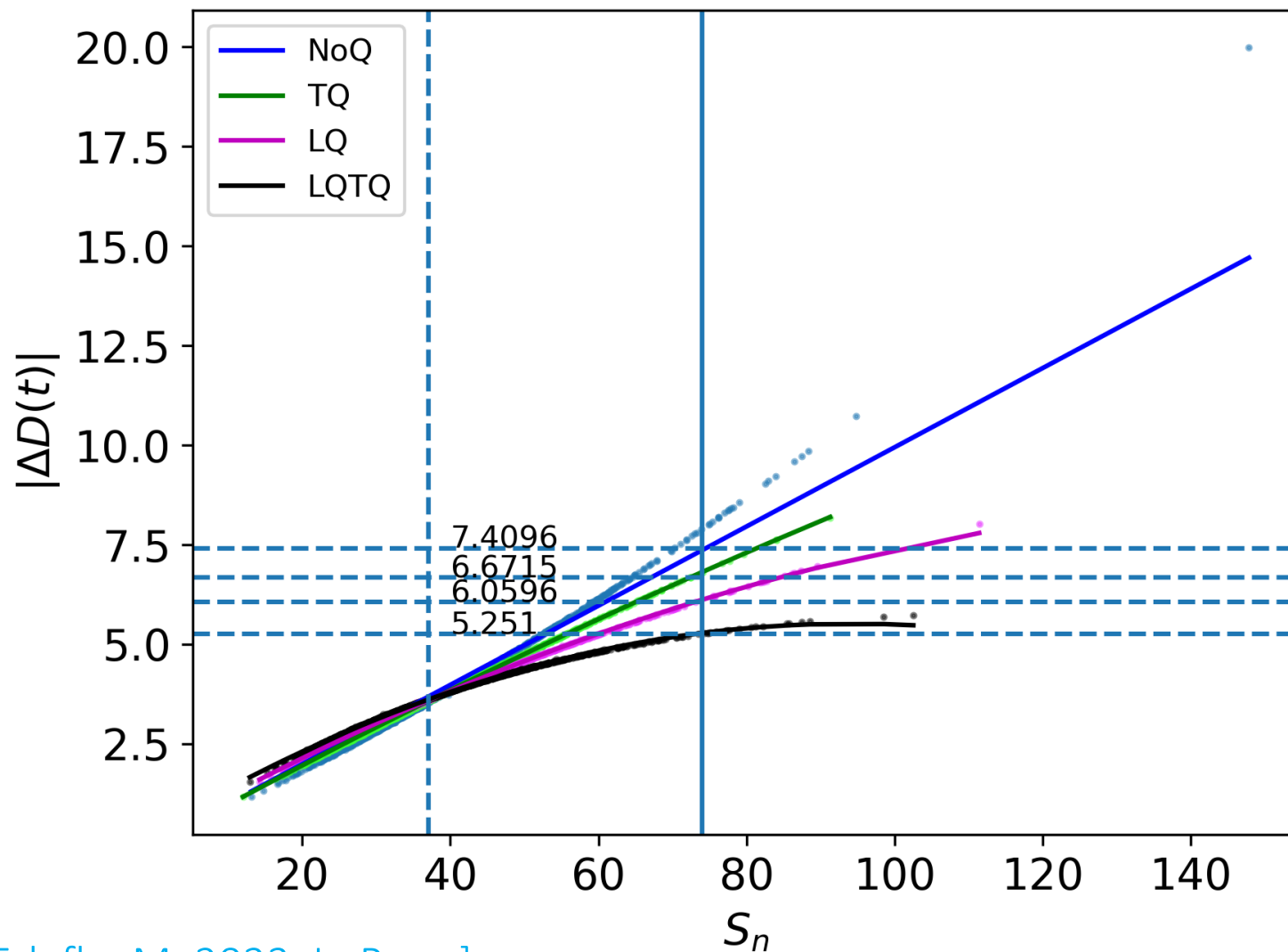
- Net contribution of a solar dipolar moment vs. cycle amplitude for the linear case and for the cases with tilt quenching and/or latitude quenching with quadratic fits. Parameter values were  $u_0 = 10\text{m/s}$ ,  $\eta = 250\text{km}^2/\text{s}$ , and  $\tau = 8$  years.
- No Inflow considered



[Talafha M. 2023, In Prep.]

$$u_0 = 10.0, \eta = 250, \tau = 8, v_0 = 5, \lambda_c = 15$$

- Net contribution of a solar dipolar moment vs. cycle amplitude for the linear case and for the cases with tilt quenching and/or latitude quenching with quadratic fits. Parameter values were  $u_0 = 10\text{m/s}$ ,  $\eta = 250\text{km}^2/\text{s}$ , and  $\tau = 8$  years.
- Inflow considered with  $v_0 = 5$  and  $\lambda_c = 15$



[Talafha M. 2023, In Prep.]

- Introducing surface inflows to the model produce lower net contribution value of solar dipolar moment -- In Agreement with the 2x2D dynamo model [Nagy, M. et. al.(2020). JSWSC, 10, 62.]

Difference between the no inflow case and the case with inflow  $v_0 = 5$  and  $\lambda_c = \pm 15$  in deviations of the nonlinear models  $|\Delta D_n|$  from the linear case measured at a certain SSN value for  $u_0 = 10$  m/s.

$\eta$ [km <sup>2</sup> /s]	$\tau$ [yr]	$\Delta dev_{LQTQ}$	$\Delta dev_{LQ}$	$\Delta dev_{TQ}$	$\Delta dev_{NoQ}$
250	8	1.4407	1.8237	1.5364	2.1424
250	$\infty$	1.7566	2.1436	1.8074	2.1273
600	8	1.2398	1.3772	1.7802	2.2745
600	$\infty$	1.5271	1.4997	2.0958	1.9429

[Talafha M. 2023, In Prep.]

# Acknowledgements

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# Thank you



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